# Solution of Transient Pressure in Radial Composite Oil Reservoir by Variable Separation Method 

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#### Abstract

Due to original sedimentary environments and accumulation processes, oil and gas reservoir are usually of great in-homogeneity. After the later-stage reservoir stimulation measures, the formation property of near-well zone and distant-well zone are often quite different. In order to characterize this kind of reservoir more accurately, the reservoir can be radially treated as composite reservoir with different zones. Based on method of variables separation to solve conventional flow problems, the pressure solution is originally obtained for radially dual composite reservoir by using the same method. This case enlarged the application of variables separation method for solving flow problems in oil and gas reservoir, and these work is also of great significance for evaluation of composite reservoirs.


Keywords: radial flow, transient flow, method of variables separation, analytical solution, composite reservoir

## 1. Introduction

With consistent focus on sustainable development for oil and gas resources in recent years, many unconventional oil and gas resources become the exploitation target for oil companies[1-3]. In order to solve these emerging seepage problems which are becoming more and more complicated in the development of unconventional oil and gas resources, it is necessary to strengthen the study on solution methods for seepage problems. Until now, many scholars have carried out detailed investigation from the perspective of analytical solution, approximate solution and numerical solution for these seepage problems, and significant progress have been made[4-7]. Compared with the practicality of numerical solution, the results obtained form analytical solution can be a good tool for sensitive analysis of each related parameters in the seepage process, and sometimes, the introduction of analytical methods can be helpful for the reduction of computing jobs. In the following, the pressure of unsteady flow in radial composite reservoir is studied, the mathematical model is established based on the characteristics of seepage, and the accurate analytical solution is finally obtained using the variable separation method.

## 2. Overview of Variable Separation Method

As for the idea, method, and procedure of variable separation method, many references have demonstrated very thoroughly in the view of wave equation[8-10]. Since the equation for seepage problems is usually parabolic equation, the case study is for solution of parabolic equation here first, and this can demonstrate effectiveness of variable separation method in solving simple seepage problems which is the foundation for solution of following seepage problems in composite reservoir. Suppose the pressure $p$ in one dimensional unsteady flow can meet the following mixed problem:

$$
\left\{\begin{array}{l}
\frac{\partial^{2} p}{\partial x^{2}}=\frac{1}{\eta} \frac{\partial p}{\partial t}, 0<x<l, t>0  \tag{1}\\
p(x, 0)=\varphi(x), 0 \leq x \leq l \\
p(0, t)=0, t \geq 0 \\
\frac{\partial p}{\partial x}+\left.h p\right|_{x=l}=0, t \geq 0
\end{array}\right.
$$

Suppose the solution can be expressed in the format of variable separation as below:

$$
\begin{equation*}
p(x, t)=X(x) T(t) \tag{2}
\end{equation*}
$$

Substitute formula(2) into constitutive equation in problem(1):

$$
\begin{align*}
T^{\prime}(t)+\eta \lambda T(t) & =0  \tag{3}\\
X^{\prime \prime}(x)+\lambda X(x) & =0 \tag{4}
\end{align*}
$$

Take formula(4) into consideration, in combination with the boundary conditions in problem(1), then function $X(x)$ can meet the following boundary condition:

$$
\begin{equation*}
X(0)=0, X^{\prime}(l)+h X(l)=0 \tag{5}
\end{equation*}
$$

For eigenvalue problem(4) and (5), it is not hard to find that it has infinitely many eigenvalues $\lambda_{k}$, $k=1,2, \cdots$. Among them, $\lambda_{k}$ is the positive roots of following equation for parameter $\lambda$ :

$$
\begin{equation*}
\sqrt{\lambda} \cos \sqrt{\lambda} l+h \sin \sqrt{\lambda} l=0 \tag{6}
\end{equation*}
$$

The corresponding eigenfunction is:

$$
\begin{equation*}
X_{k}(x)=B_{k} \sin \sqrt{\lambda_{k}} x, k=1,2, \cdots \tag{7}
\end{equation*}
$$

Substitute $\lambda_{k}$ into formula(3), there will be:

$$
T_{k}(t)=A_{k} e^{-\eta \lambda_{k} t}, k=1,2, \cdots \text { (8) }
$$

Therefore, suppose $A_{k} B_{k}=C_{k}$, a series of special solutions with variable separation format can be obtained:

$$
p_{k}(x, t)=C_{k} e^{-\eta \lambda_{k} t} \sin \sqrt{\lambda_{k}} x, k=1,2, \cdots(9)
$$

Consider the linearity and homogeneity of equation and boundary conditions in problem(1), the solution expressed in series can be obtained based on the theory of superposition:

$$
\begin{equation*}
p(x, t)=\sum_{k=1}^{\infty} C_{k} e^{-\eta \lambda_{k} t} \sin \sqrt{\lambda_{k}} x \tag{10}
\end{equation*}
$$

In combination with the initial condition in problem(1), expand function $\varphi(x)$ into sine series, then their coefficients $C_{k}$ can be determined by the following formula:

$$
\begin{equation*}
C_{k}=\frac{1}{M_{k}} \int_{0}^{l} \varphi(\xi) \sin \sqrt{\lambda_{k}} \xi d \xi \tag{11}
\end{equation*}
$$

Among them, there is:

$$
\begin{equation*}
M_{k}=\int_{0}^{l} \sin ^{2} \sqrt{\lambda_{k}} \xi d \xi \tag{12}
\end{equation*}
$$

Substitute formula(11)-(12) into formula(10), and the solution of problem 1) in variable separation format will finally be obtained.

## 3. Mathematical Model Establishment for Seepage Problem

Consider to enlarge the application of above variable separation method, suppose the radial twoegion composite reservoir is as that shown in figure 1: the formation is horizontal and it consists of two regions with permeability of $K_{1}$ and $K_{2}$ respectively. The two regions are distributed radially and the interval ranges are $0 \leq r \leq a$ and $a \leq r \leq b$ respectively, and the pressure transmission coefficients
are $\eta_{1}$ and $\eta_{2}$ respectively. The pressure at the outer boundary remains constant. After production for some time, the pressure change from that(the initial formation pressure is $p_{i}$ ) at the initial time $t=0$ are $F_{1}(r)$ and $F_{2}(r)$ respectively. Suppose there are good connectivity between two regions, and the purpose is to find out the pressure distribution in each region when $t>0$.


Figure 1: Sketch map of seepage area for radial two-region composite reservoir.
In order to make outer boundary conditions homogeneous, suppose $p_{1}(r, t)$ and $p_{2}(r, t)$ are the pressure change compared with that of initial pressure in two regions respectively, refer to the model establishment methods for seepage problem[11-12], the mathematical expression of this problem can be established as:

$$
\begin{array}{r}
\frac{1}{r} \frac{\partial}{\partial r}\left(r \frac{\partial p_{1}}{\partial r}\right)=\frac{1}{\eta_{1}} \frac{\partial p_{1}}{\partial t}, \quad 0<r<a(13) \\
\frac{1}{r} \frac{\partial}{\partial r}\left(r \frac{\partial p_{2}}{\partial r}\right)=\frac{1}{\eta_{2}} \frac{\partial p_{2}}{\partial t}, \quad a<r<b(14) \\
p_{1}=p_{2}, \quad r=a(15) \\
K_{1} \frac{\partial p_{1}}{\partial r}=K_{2} \frac{\partial p_{2}}{\partial r}, \quad r=a(16) \\
p_{2}=0, \quad r=b(17) \\
p_{1}(r, t)=F_{1}(r), \quad 0 \leq r \leq a, \quad t=0(18) \\
p_{2}(r, t)=F_{2}(r), \quad a \leq r \leq b, \quad t=0(19) \tag{19}
\end{array}
$$

## 4. Solution of Seepage Problem by Variable Separation Method

Suppose $p_{1}(r, t)=R_{1}(r) T_{1}(t), p_{2}(r, t)=R_{2}(r) T_{2}(t)$, substitute into the formulas(13)-(19), introduce the separation constant $\lambda \equiv \beta^{2}$, therefore, the separation results will be:

$$
\begin{equation*}
T^{\prime}+\beta^{2} T=0 \tag{20}
\end{equation*}
$$

and $R_{1}(\beta, r), R_{2}(\beta, r)$ can meet the following eigenvalue problems:

$$
\left\{\begin{array}{l}
\frac{1}{r} \frac{d}{d r}\left(r \frac{d R_{1}}{d r}\right)+\frac{\beta^{2}}{\eta_{1}} R_{1}=0,0<r<a  \tag{21}\\
\frac{1}{r} \frac{d}{d r}\left(r \frac{d R_{2}}{d r}\right)+\frac{\beta^{2}}{\eta_{2}} R_{2}=0, a<r<b \\
R_{1}=R_{2}, r=a \\
K_{1} \frac{d R_{1}}{d r}=K_{2} \frac{d R_{2}}{d r}, r=a \\
R_{2}=0, r=b
\end{array}\right.
$$

Based on formula(20), the following results can be obtained:

$$
\begin{equation*}
T(t)=e^{-\beta^{2} t} \tag{22}
\end{equation*}
$$

Problem(21) is the eigenvalue problem of this composite reservoir, the eigenvalue $\beta_{m}$ and the eigenfunction can be solved and expressed in two intervals:

$$
R\left(\beta_{m}, r\right)=\left\{\begin{array}{l}
R_{1}\left(\beta_{m}, r\right), 0 \leq r \leq a  \tag{23}\\
R_{2}\left(\beta_{m}, r\right), a \leq r \leq b
\end{array}\right.
$$

It can be proved that the eigenfunction has following orthogonality:

$$
\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}\left(\beta_{m}, r\right) R_{1}\left(\beta_{n}, r\right) d r+\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}\left(\beta_{m}, r\right) R_{2}\left(\beta_{n}, r\right) d r=\left\{\begin{array}{l}
0, \quad m \neq n  \tag{24}\\
N\left(\beta_{m}\right), \quad m=n
\end{array}\right.
$$

Among them, the norm is:

$$
\begin{equation*}
N\left(\beta_{m}\right)=\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}^{2}\left(\beta_{m}, r\right) d r+\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}^{2}\left(\beta_{m}, r\right) d r \tag{25}
\end{equation*}
$$

So the solution of problem(13)-(19) can be expressed in two intervals:

$$
p(r, t)=\left\{\begin{array}{l}
p_{1}(r, t), 0 \leq r \leq a  \tag{26}\\
p_{2}(r, t), a \leq r \leq b
\end{array}\right.
$$

Among them,

$$
\begin{align*}
& p_{1}(r, t)=\sum_{m=1}^{\infty} c_{m} e^{-\beta_{m}^{2} t} R_{1}\left(\beta_{m}, r\right), 0 \leq r \leq a  \tag{27}\\
& p_{2}(r, t)=\sum_{m=1}^{\infty} c_{m} e^{-\beta_{m}^{2} t} R_{2}\left(\beta_{m}, r\right), a \leq r \leq b \tag{28}
\end{align*}
$$

Take initial condition into consideration, the superposition coefficient $c_{m}$ can be determined:

$$
\begin{align*}
& F_{1}(r)=\sum_{m=1}^{\infty} c_{m} R_{1}\left(\beta_{m}, r\right), 0 \leq r \leq a  \tag{29}\\
& F_{2}(r)=\sum_{m=1}^{\infty} c_{m} R_{2}\left(\beta_{m}, r\right), a \leq r \leq b \tag{30}
\end{align*}
$$

Make use of orthogonality formula(24), do operation on each side of formulas (29) and (30) with
operator $\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}\left(\beta_{m}, r\right) d r$ and $\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}\left(\beta_{m}, r\right) d r$, and then add results together, there will be:

$$
\begin{equation*}
c_{m}=\frac{1}{N\left(\beta_{m}\right)}\left[\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}\left(\beta_{m}, r\right) F_{1}(r) d r+\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}\left(\beta_{m}, r\right) F_{2}(r) d r\right] \tag{31}
\end{equation*}
$$

Substitute formula(31) into formula(27) and (28), there will be:

$$
\begin{align*}
& p_{1}(r, t)=\sum_{m=1}^{\infty} e^{-\beta_{m}{ }^{2} t} R_{1}\left(\beta_{m}, r\right) \frac{1}{N\left(\beta_{m}\right)} .  \tag{32}\\
& {\left[\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}\left(\beta_{m}, r\right) F_{1}(r) d r+\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}\left(\beta_{m}, r\right) F_{2}(r) d r\right]} \\
& p_{2}(r, t)=\sum_{m=1}^{\infty} e^{-\beta_{m}^{2} t} R_{2}\left(\beta_{m}, r\right) \frac{1}{N\left(\beta_{m}\right)} \cdot  \tag{33}\\
& {\left[\frac{K_{1}}{\eta_{1}} \int_{0}^{a} r R_{1}\left(\beta_{m}, r\right) F_{1}(r) d r+\frac{K_{2}}{\eta_{2}} \int_{a}^{b} r R_{2}\left(\beta_{m}, r\right) F_{2}(r) d r\right]}
\end{align*}
$$

By now, the formal solution has been obtained, and the remaining task is to solve the eigenvalue problem 21):

First, the general solution of equations 21) can be solved as:

$$
\begin{align*}
& R_{1}(\beta, r)=A_{1} J_{0}\left(\frac{\beta}{\sqrt{\eta_{1}}} r\right)+B_{1} N_{0}\left(\frac{\beta}{\sqrt{\eta_{1}}} r\right), 0 \leq r \leq a  \tag{34}\\
& R_{2}(\beta, r)=A_{2} J_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} r\right)+B_{2} N_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} r\right), a \leq r \leq b \tag{35}
\end{align*}
$$

The unknown coefficients $A_{1}, B_{1}, A_{2}, B_{2}$ and the eigenvalue $\beta$ has to be determined by boundary conditions and connecting conditions.

Take the engineering reality into consideration, $p_{1}$ is certain to have limitation at $r=0$, so $R_{1}$ also has limitation correspondingly. Without loss of generality, suppose $A_{1}=1$, formula (34) and (35) above can be simply expressed as:

$$
\begin{align*}
& \quad R_{1}(\beta, r)=J_{0}\left(\frac{\beta}{\sqrt{\eta_{1}}} r\right), 0 \leq r \leq a  \tag{36}\\
& R_{2}(\beta, r)=A_{2} J_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} r\right)+B_{2} N_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} r\right),  \tag{37}\\
& a \leq r \leq b
\end{align*}
$$

Make use of the boundary condition at $r=b$ and the connecting condition at $r=a$, the equations concerning $A_{2}, B_{2}$, and $\beta$ can be obtained as follows:

$$
\begin{array}{r}
J_{0}\left(\frac{\beta}{\sqrt{\eta_{1}}} a\right)=A_{2} J_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} a\right)+B_{2} N_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} a\right) \\
\frac{K_{1}}{K_{2}} \sqrt{\frac{\eta_{2}}{\eta_{1}} J_{1}\left(\frac{\beta}{\sqrt{\eta_{1}}} a\right)=A_{2} J_{1}\left(\frac{\beta}{\sqrt{\eta_{2}}} a\right)+B_{2} N_{1}\left(\frac{\beta}{\sqrt{\eta_{2}}} a\right)} \tag{39}
\end{array}
$$

$$
\begin{equation*}
0=A_{2} J_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} b\right)+B_{2} N_{0}\left(\frac{\beta}{\sqrt{\eta_{2}}} b\right) \tag{40}
\end{equation*}
$$

Or formulas(38)-(40) can be expressed in following equation of matrix:

$$
\left[\begin{array}{ccc}
J_{0}\left(\xi_{1}\right) & -J_{0}\left(\frac{a}{b} \xi_{2}\right) & -N_{0}\left(\frac{a}{b} \xi_{2}\right) \\
K J_{1}\left(\xi_{1}\right) & -J_{1}\left(\frac{a}{b} \xi_{2}\right) & -N_{1}\left(\frac{a}{b} \xi_{2}\right) \\
0 & J_{0}\left(\xi_{2}\right) & N_{0}\left(\xi_{2}\right)
\end{array}\right] \cdot\left[\begin{array}{c}
1 \\
A_{2} \\
B_{2}
\end{array}\right]=\left[\begin{array}{l}
0 \\
0 \\
0
\end{array}\right]
$$

Among them,

$$
\xi_{1}=\frac{\beta}{\sqrt{\eta_{1}}} a, \xi_{2}=\frac{\beta}{\sqrt{\eta_{2}}} b, K=\frac{K_{1}}{K_{2}} \sqrt{\frac{\eta_{2}}{\eta_{1}}}
$$

Since the equation above has non-zero roots, the determinant of its coefficients must be 0 , there is:

$$
\left|\begin{array}{ccc}
J_{0}\left(\xi_{1}\right) & -J_{0}\left(\frac{a}{b} \xi_{2}\right) & -N_{0}\left(\frac{a}{b} \xi_{2}\right) \\
K J_{1}\left(\xi_{1}\right) & -J_{1}\left(\frac{a}{b} \xi_{2}\right) & -N_{1}\left(\frac{a}{b} \xi_{2}\right) \\
0 & J_{0}\left(\xi_{2}\right) & N_{0}\left(\xi_{2}\right)
\end{array}\right|=0
$$

Eigenvalue $\beta$ is just the root of this equation. There are infinitely many roots: $\beta=\beta_{m}, m=1,2, \cdots$.

Due to the validity of formula(41), formula(38)-(40) are not totally independent, therefore, any two of the three formulas can be chosen for determination of coefficients. Coefficients $A_{2}$ and $B_{2}$ can be determined by combination of formula(38) and (40):

$$
\begin{align*}
A_{2} & =\frac{J_{0}\left(\xi_{1}\right) N_{0}\left(\xi_{2}\right)}{M}  \tag{42}\\
B_{2} & =-\frac{J_{0}\left(\xi_{1}\right) J_{0}\left(\xi_{2}\right)}{M} \tag{43}
\end{align*}
$$

Among them,

$$
\mathrm{M}=\left|\begin{array}{cc}
J_{0}\left(\frac{a}{b} \xi_{2}\right) & N_{0}\left(\frac{a}{b} \xi_{2}\right)  \tag{44}\\
J_{0}\left(\xi_{2}\right) & N_{0}\left(\xi_{2}\right)
\end{array}\right|
$$

Substitute the results of formula(42)-(44) into formula(37), the eigenfunctions can be completely determined. Actually, the coefficients $A_{2}, B_{2}$ are actually related to eigenvalue, for each eigenvalue $\beta_{m}$, there are corresponding coefficients $A_{2 m}, B_{2 m}$.

## 5. Conclusions

The solutions directly obtained is actually the pressure change compared with that of original formation pressure. To get the pressure distribution in each region, just add the initial pressure $p_{i}$ with $p_{1}(r, t)$ or $p_{2}(r, t)$ respectively.

Choose the seepage problem in radial composite reservoir as an example, the joining condition between two regions are considered, on the basis of solution of seepage problems by traditional variable separation method, the solution of pressure in composite reservoir is obtained. This study enlarged the application of variable separation method, and which significantly improve the methodology for the evaluation of oil and gas reservoirs.

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## References

[1] Jiao Fangzheng, (2019) Re-recognition of "unconventional" in unconventional oil and gas. Petroleum Exploration and Development,46(5):803-810.
[2] Zou Caineng, Yang Zhi, Zhang Guosheng, et al. (2019) Establishment and practice of unconventional oil and gas geology. Acta Geologica Sinica,93(1):12-23.
[3] Guo Jin, (2019) Characteristics and Classification of Unconventional Oil and Gas Reservoirs. Petrochemical Industry Technology,26(6):135-136,141.
[4] Li Zhiqiang, and Cheng Qiang, (2018) Multi-process seepage engineering application and its numerical solutions in gas extraction field. Coal Engineering, 50(11):74-78.
[5] Xue Dongjie,Zhou Hongwei,Deng Linsheng, et al. (2018) Fractal Dynamics of Gas-liquid Flow in Low-permeability Coal. Journal of Sichuan University(Engineering Science Edition),50(4):30-40.
[6] Zhang Liehui,Liu Xiangyu,Zhao Yulong, et al. (2019) Effect of pore throat structure on micro-scale seepage characteristics of tight gas reservoirs. Natural Gas Industry, 39(8):50-57.
[7] Hou Shaoji, Zhu Weiping, Liu Yuewu, et al. (2018) A Semi-Analytical Model for Moving Boundary of Radial Non-Darcy Flow in Low Permeability Reservoir.Applied Mathematics and Mechanics, 39(10):1115-1127.
[8] Yao Duanzheng, Liang Jiabao, (2011) Methods of mathematical physics, 2nd edn. Wuhan University Press, Wuhan.
[9] Ke Daoming, Huang Zhixiang, Chen Junning, (2018) Methods of mathematical physics, 2nd edn. China Machine Press, Beijing.
[10] Wang Mingxin, (2009) Equations of mathematical physics, 2nd edn. China Machine Press, Beijing. [11] Kong Xiangyan, (2010) Advanced Flow in Porous Media, 2nd edn. Press of University of Science and Technology of China, Hefei.
[12] Chen Junbin, Wang Bing, Zhang Guoqiang, (2013) Dynamics and physics offuilds in porous media. Petroleum Industry Press, Beijing.

